

Magnetic torque anomaly in Dirac semimetal Cd_3As_2

¹F. Orbanic, ¹M. Novak, ²A. McCollam, ²L. Tang, ¹I. Kokanovic

¹Department of Physics, Faculty of Science, University of Zagreb, Croatia

²High Field Magnet Laboratory, Radboud University, Nijmegen, the Netherlands

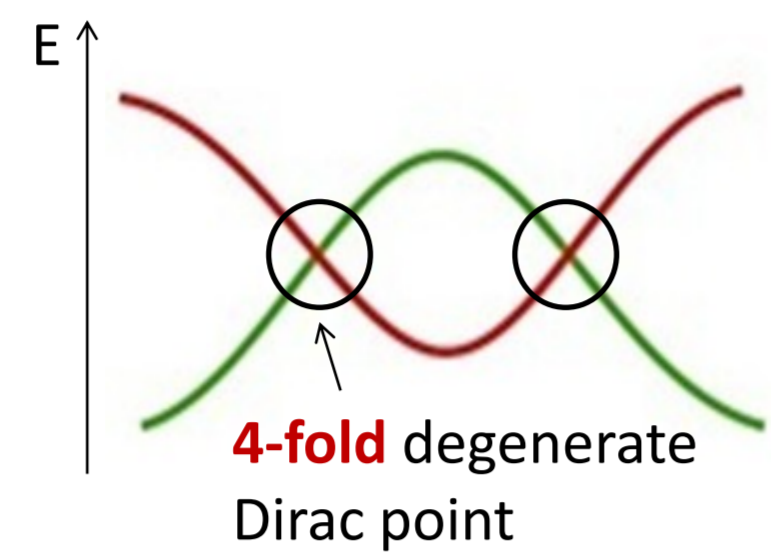
E-mail: forbanic@phy.hr

3D Dirac semimetal

Dirac dispersion in 3D (k -space) \rightarrow 3D analogue of graphene.

Dirac fermion physics:

- High mobility and low effective mass.
- Large LMR.
- Interesting transport properties.
- Fundamental physics (Dirac/Weyl).



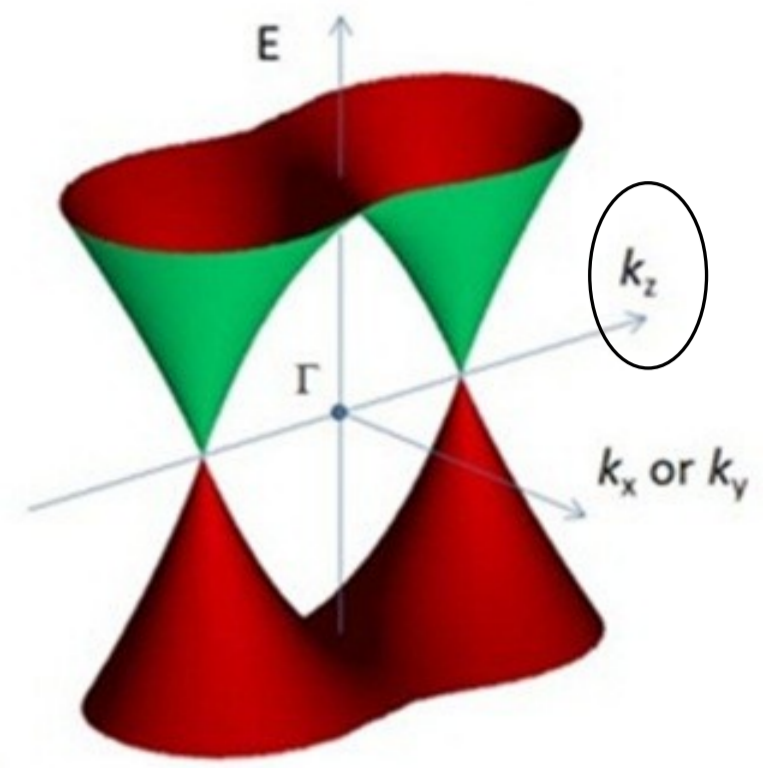
Cd_3As_2

\rightarrow 3D Dirac semimetal with symmetry protected pair of Dirac points at Z- Γ -Z line.

Single crystals synthesized by modified CVD technique.



B-field along symmetry axis can split Dirac point in two Weyl points.



Downloaded from [2]

Transport measurements

$$m_c^* = 0.028m_e$$

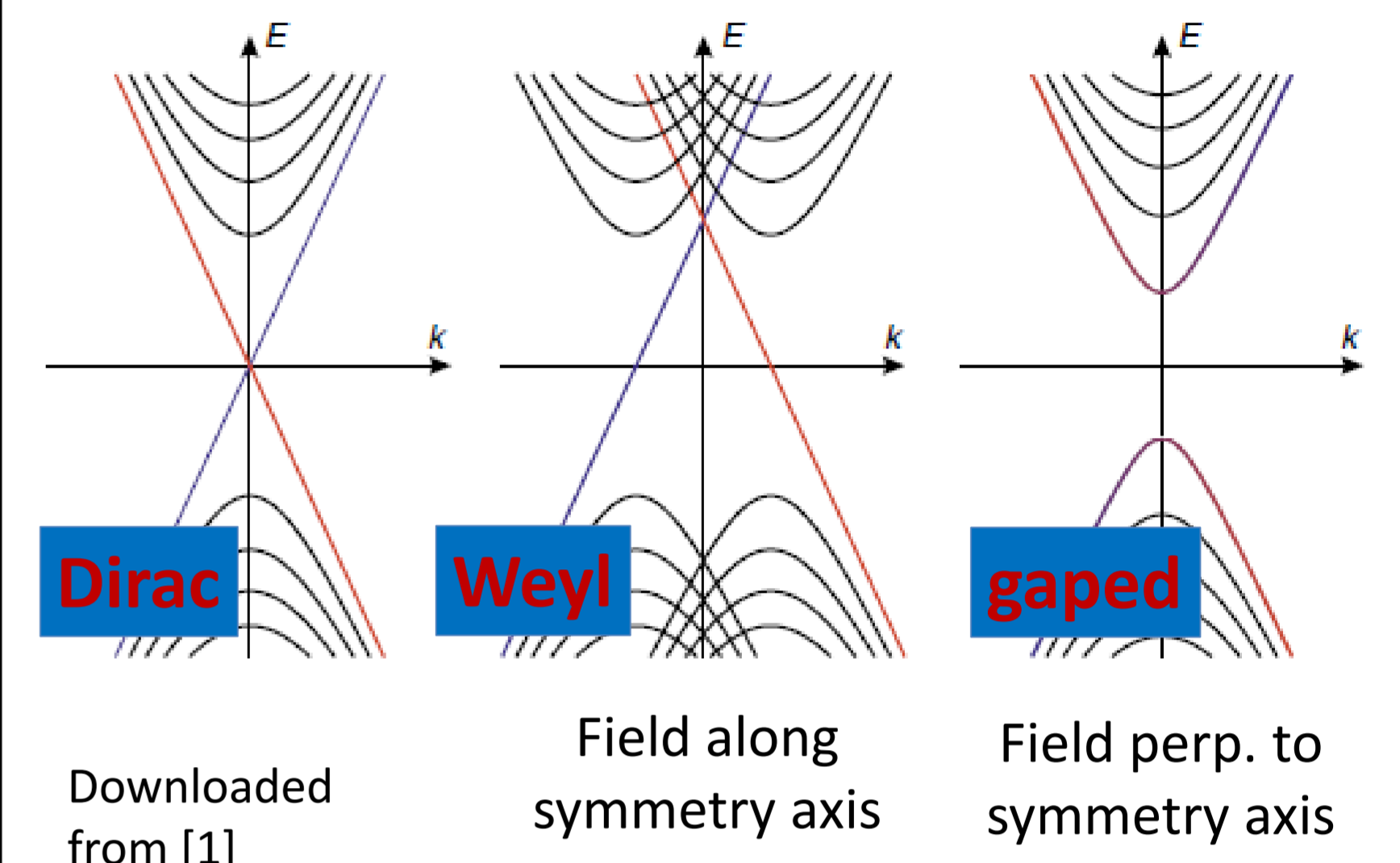
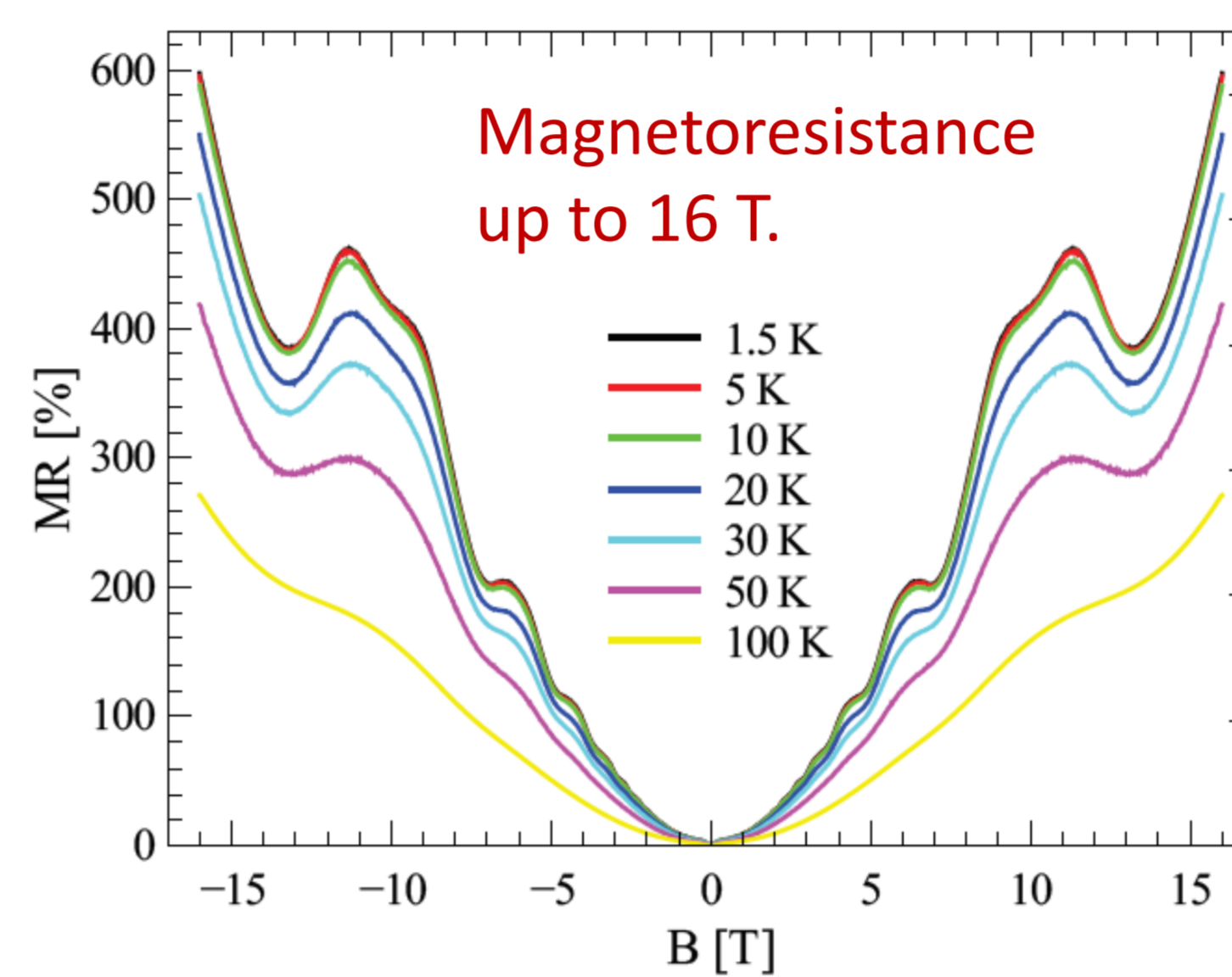
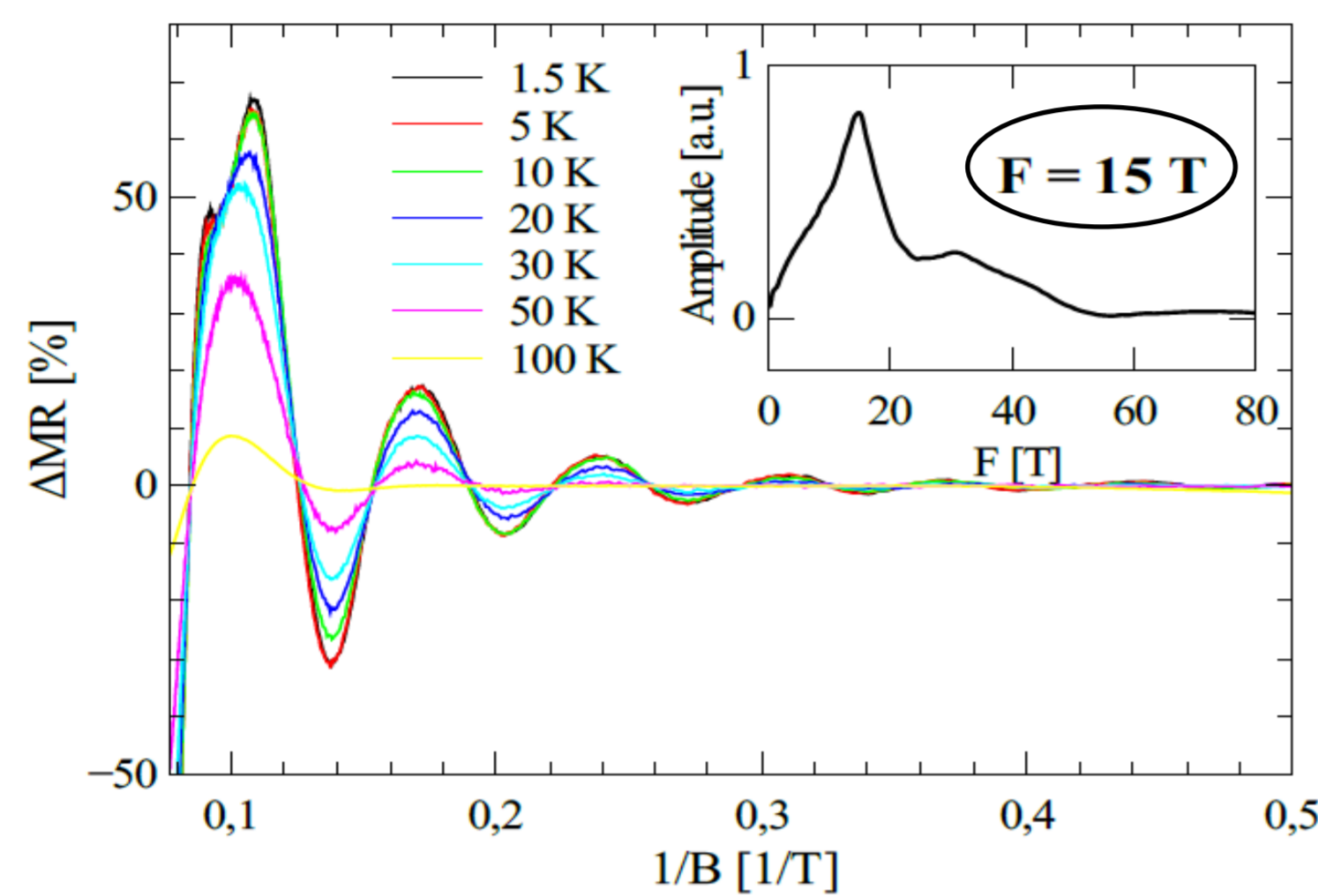
$$T_D = 35.7 \text{ K}$$

$$\tau_Q = 2.14 \cdot 10^{-13} \text{ s}$$

$$\mu = 1.34 \cdot 10^4 \text{ cm}^2/\text{Vs}$$

$$n = 3.28 \cdot 10^{17} \text{ cm}^{-3}$$

$$E_F = 124 \text{ meV}$$



Magnetization in the quantum limit

Energy levels in B-field for different types of electrons [1]:

$$\epsilon_{n,k} = \begin{cases} \frac{\hbar e B}{m} (n + \gamma) + \frac{\hbar^2 k_z^2}{2m} & \text{Trivial } (\gamma = \frac{1}{2}) \\ \hbar v_F \sqrt{2B(n + \gamma) + k_z^2} & \text{Weyl } (\gamma = 0) \\ \hbar v_F \sqrt{2B(n + \gamma + C^2 \sin^2 \theta) + k_z^2} & \text{Dirac } (\gamma = 0) \end{cases}$$

θ – angle between B and line connecting two Dirac points
C – material dependent parameter

Cause of nontrivial Berry phase. There is a zero LL at $\epsilon = 0$.

Weyl:
 $n > 0$ – paramagnetic
 $n < 0$ – diamagnetic

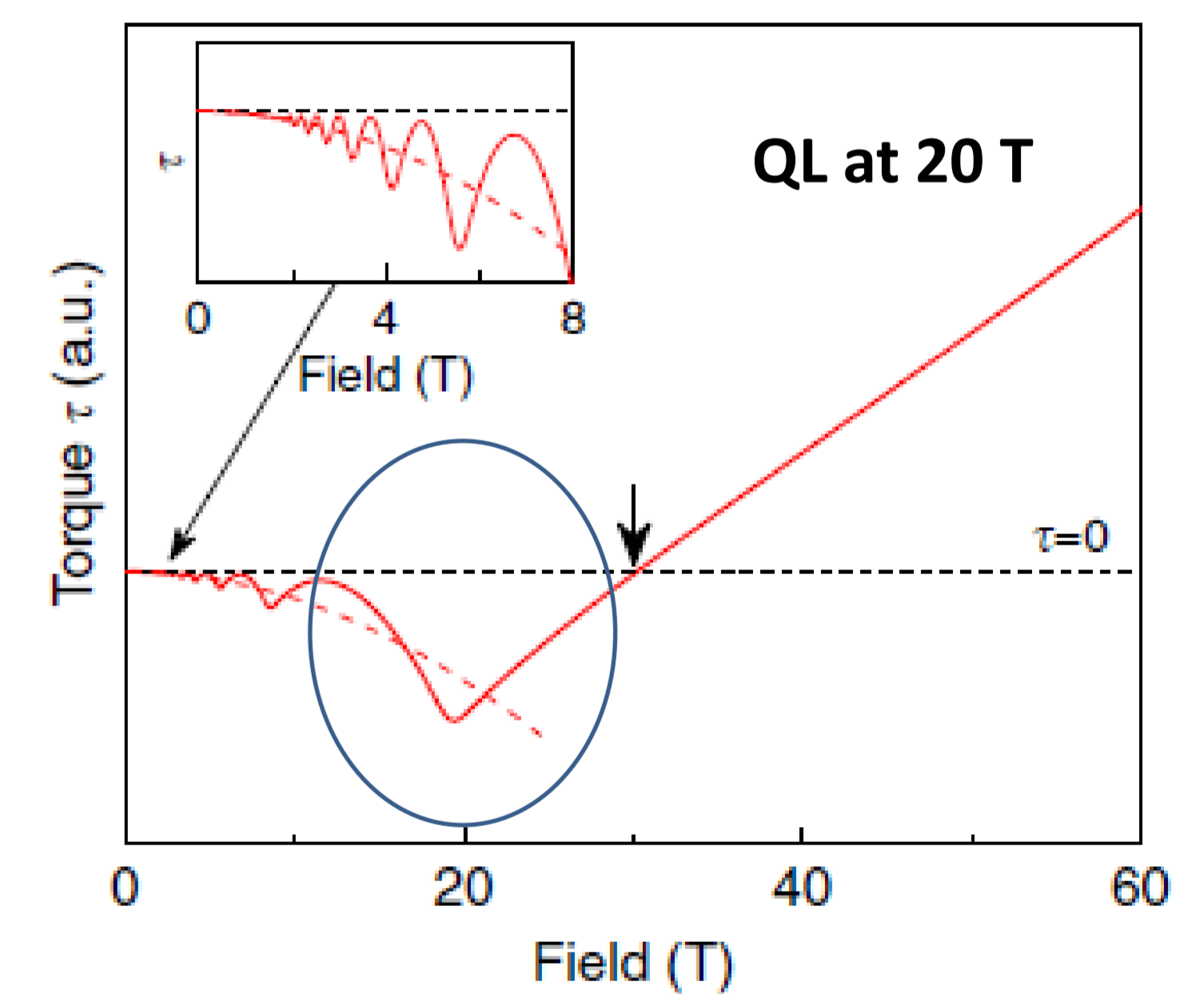
Paramagnetic response dies in QL:

$$M_{n=0} = -\frac{\partial \epsilon_{0,k}}{\partial B} = 0$$

Magnetic anomaly

Magnetic torque $T = \vec{M} \times \vec{B}$ in QL can distinguish between gaped and Weyl fermions.

Simulated behaviour of torque near the quantum limit [1]:

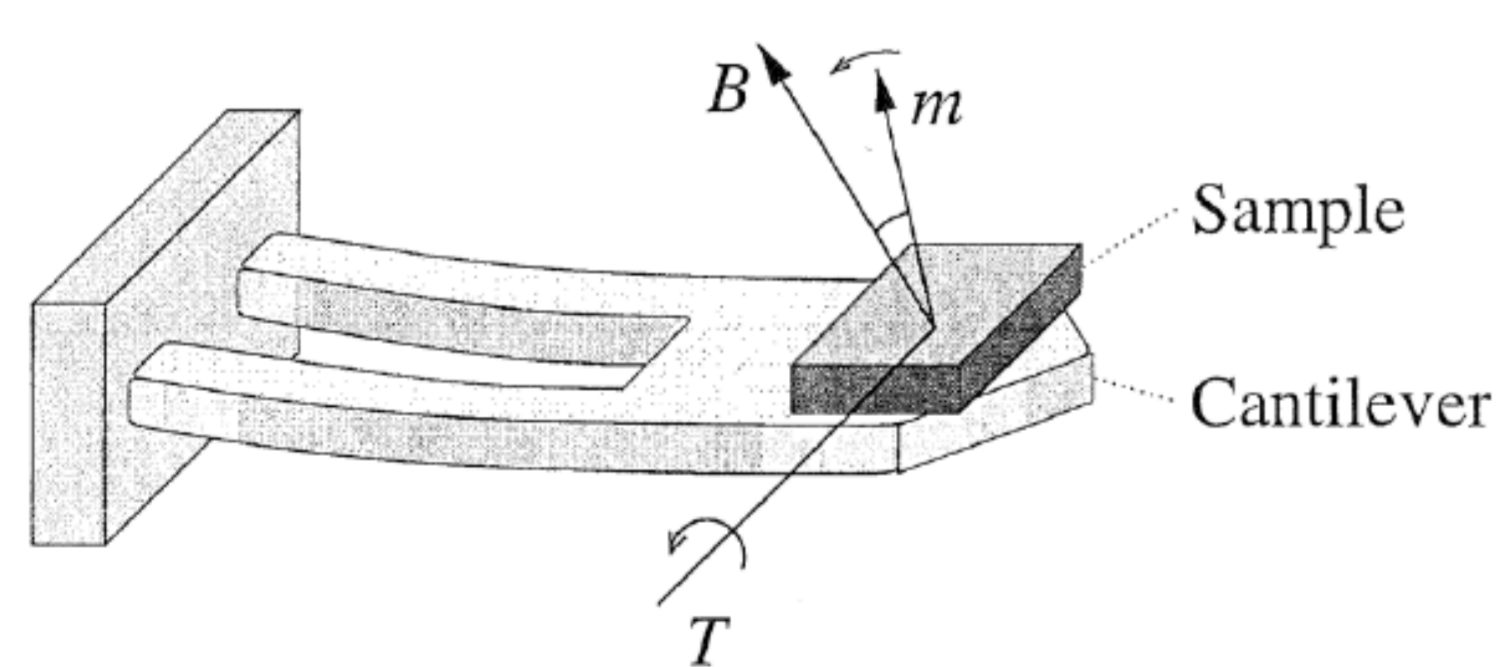


Weyl – change from diamagnetic to paramagnetic response leading to the torque reversal.

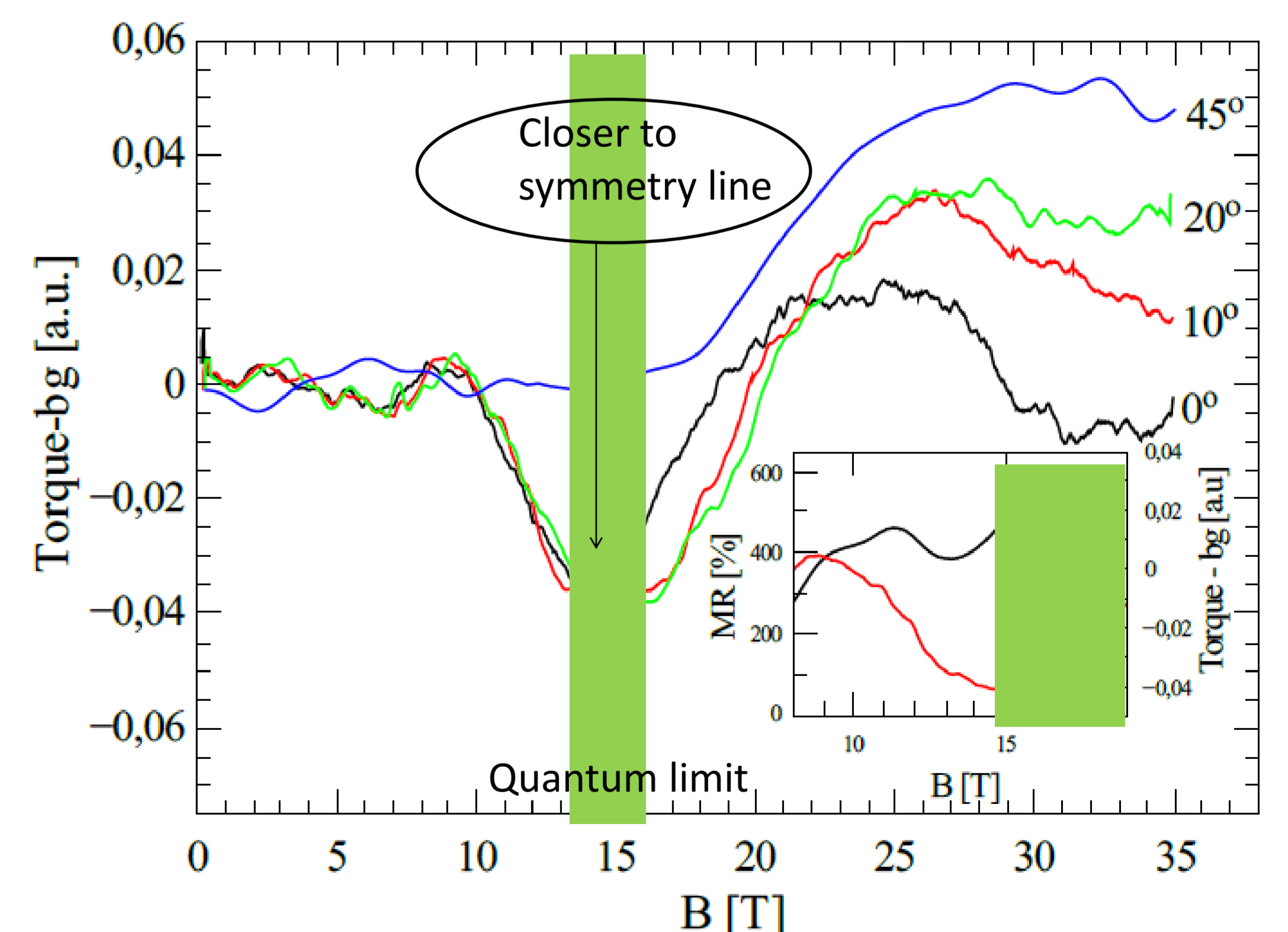
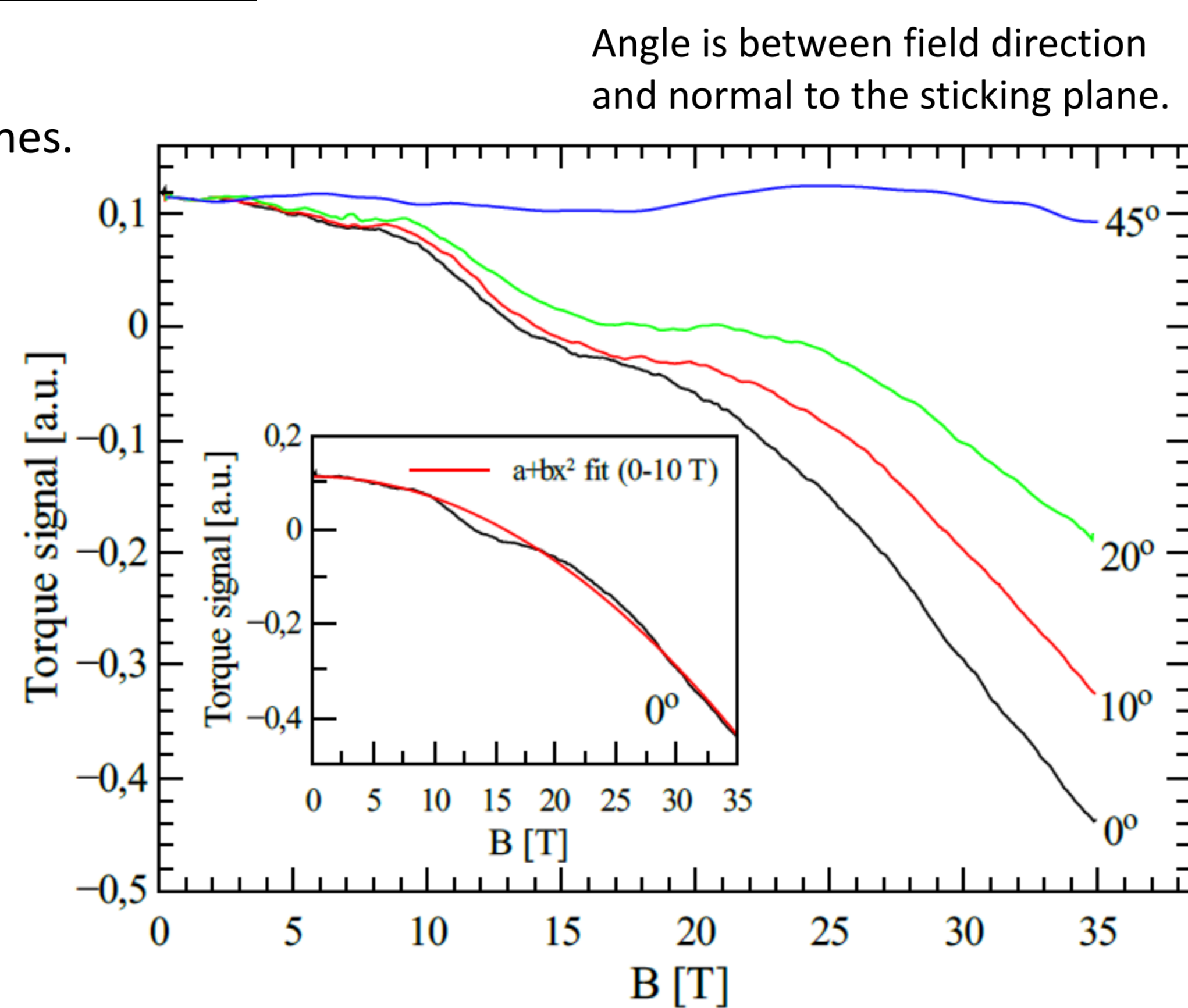
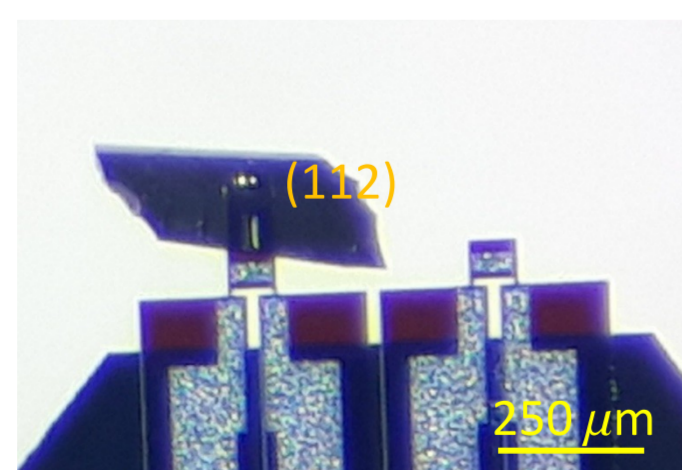
Dirac \rightarrow B-field perpendicular to the symmetry axis \rightarrow gaped semimetal
 \rightarrow B-field in the direction of the symmetry axis \rightarrow weyl

Magnetic torque measurements

- By piezoresistive cantilever technique.
- Two different samples glued on different planes.
- At 4.2 K up to 35 T.



If material possesses anisotropy, sample produces torque on lever, $T = \vec{M} \times \vec{B} \propto B^2$.



By entering the quantum limit torque anomaly appears. Anomaly and change in torque slope is highly suppressed at 45°.

Quantum limit appears at the same field for different field directions indicating the sphericity of the Fermi surface (confirmed in SdH oscillations).

Conclusion

- Successfully synthesized low charge concentration Cd_3As_2 samples with lower quantum limit (Cd_3As_2 usually has F around 45-50 T).
- SdH oscillations reveal spherical Fermi surface leading to very small torque signal. In samples with higher charge concentration the Fermi surface is ellipsoidal [2].
- Anomalous and angle dependent torque signal near the quantum limit associated to transition between Dirac and Weyl semimetal has been found.

References:

- [1] Moll, P. J. *et al.* (2016). Magnetic torque anomaly in the quantum limit of Weyl semimetals. *Nature Communications*, 7.
- [2] Borisenko, S. *et al.* (2014). Experimental realization of a three-dimensional Dirac semimetal. *Physical Review Letters*, 113(2).