

Magnetic torque in Dirac semimetal Cd_3As_2

¹F. Orbanic, ¹M. Novak, ²A. McCollam, ²L. Tang, ¹I. Kokanovic

¹Department of Physics, Faculty of Science, University of Zagreb, Croatia

²High Field Magnet Laboratory, Radboud University, Nijmegen, the Netherlands



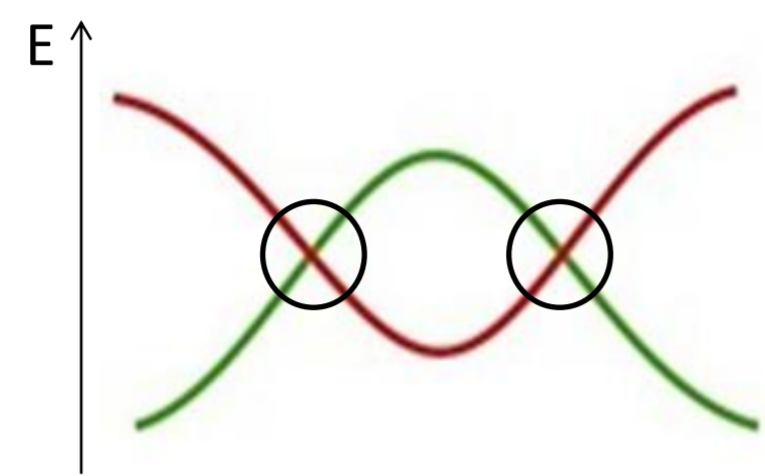
E-mail: forbanic@phy.hr

3D Dirac semimetal

Dirac dispersion in 3D (k -space) \rightarrow 3D analogues of graphene.

Dirac fermion physics:

- High mobility and low effective mass.
- Large LMR.
- Interesting transport properties.
- Fundamental physics (Weyl semimetal).



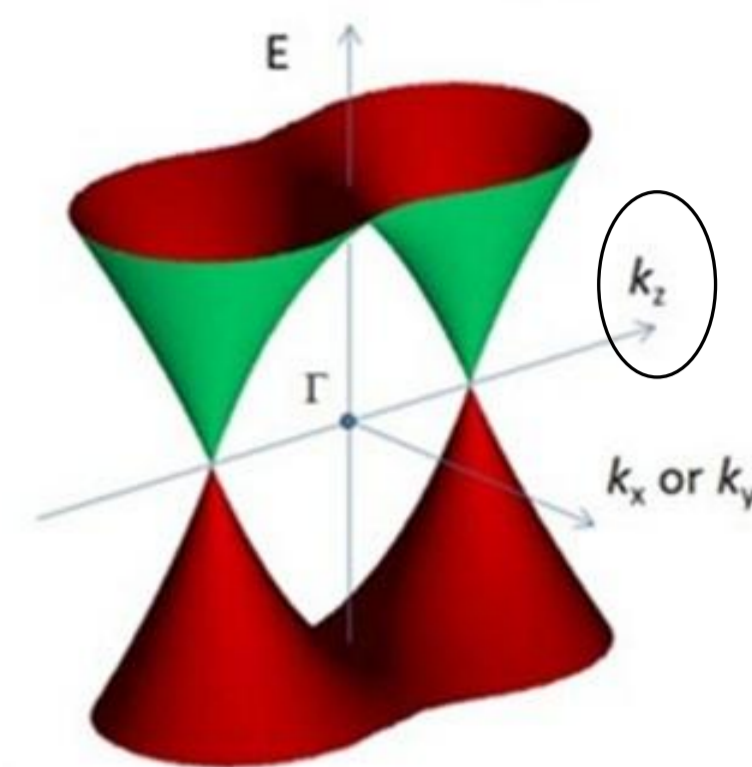
Cd_3As_2

\rightarrow 3D Dirac semimetal with symmetry protected pair of Dirac points at Z-F-Z line.

Single crystals synthesized by modified CVD technique.



B-field along symmetry axis can split Dirac point into two Weyl points.



Downloaded from [2]

Transport measurements

Magnetoresistance up to 16 T.

$$m_c^* = 0.028m_e$$

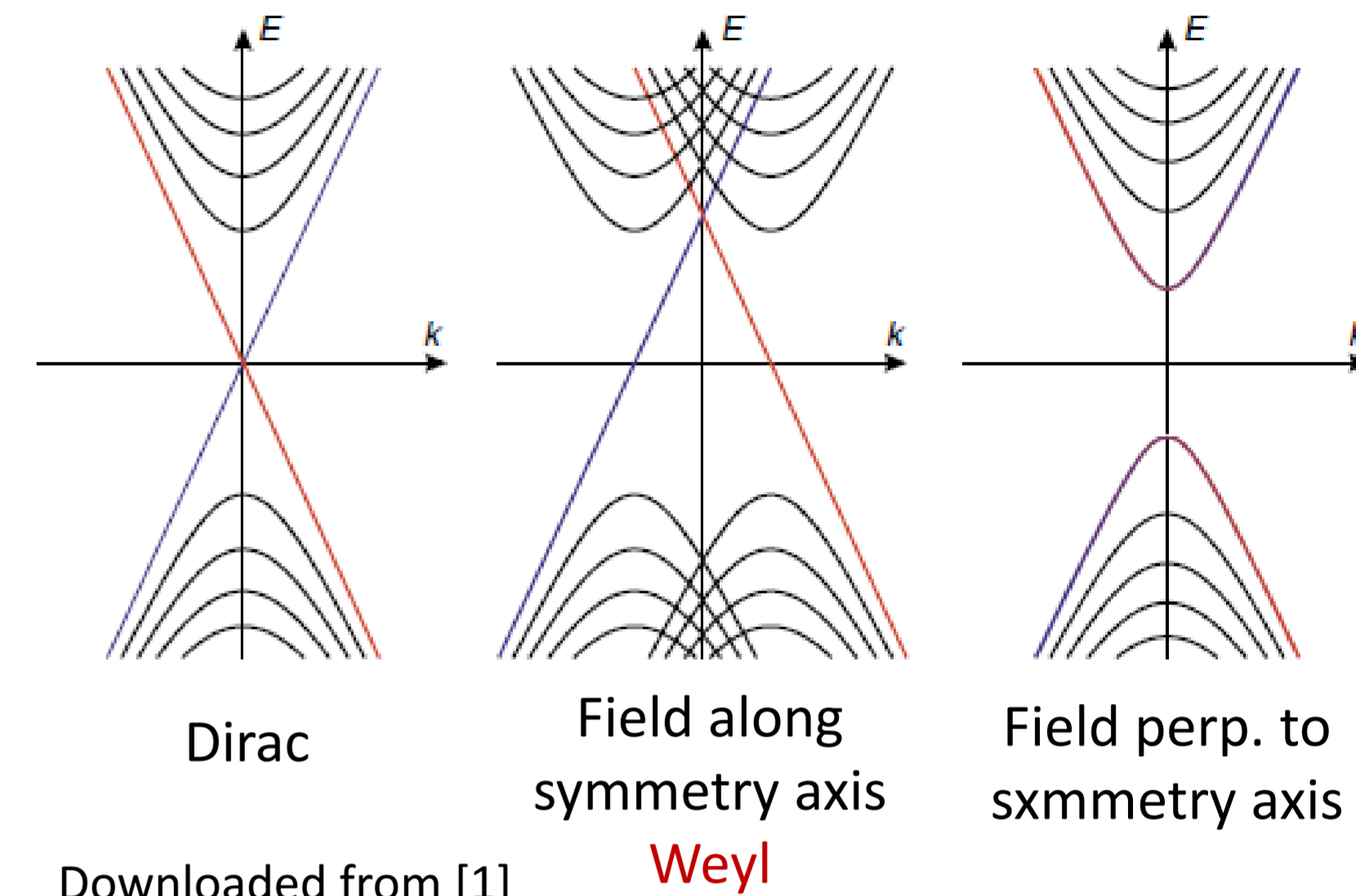
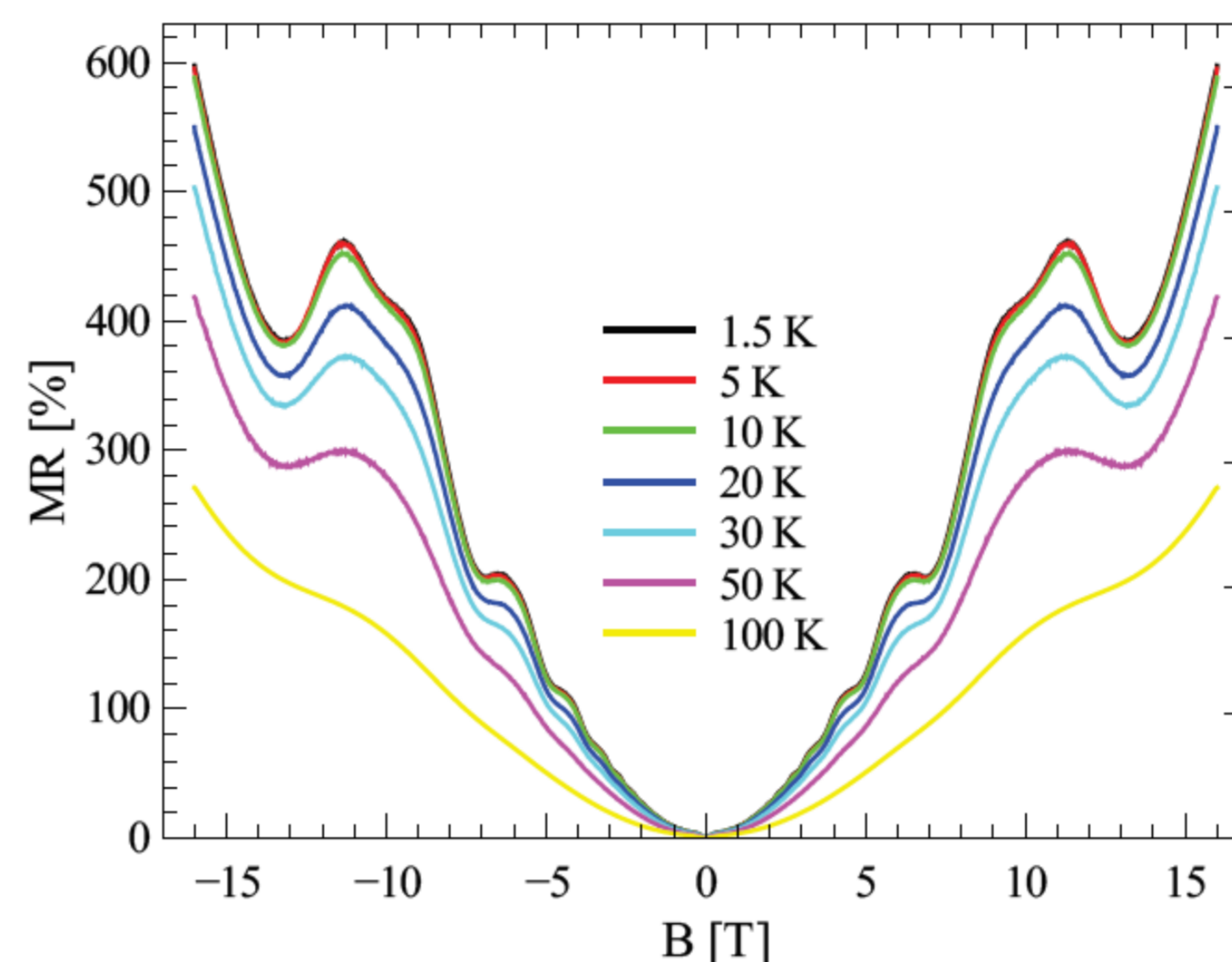
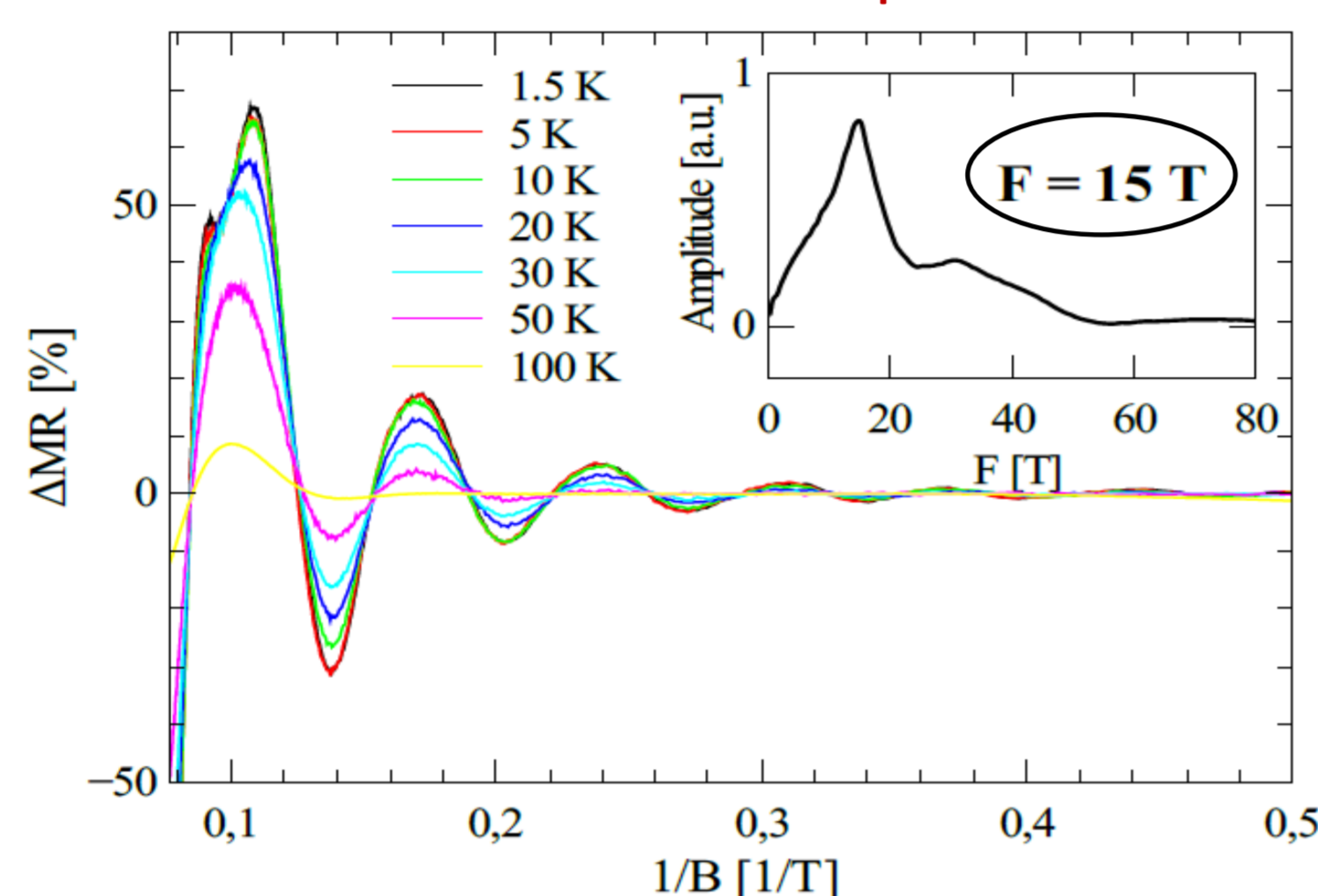
$$T_D = 35.7 \text{ K}$$

$$\tau_Q = 2.14 \cdot 10^{-13} \text{ s}$$

$$\mu = 1.34 \cdot 10^4 \text{ cm}^2/\text{Vs}$$

$$n = 3.28 \cdot 10^{17} \text{ cm}^{-3}$$

$$E_F = 124 \text{ meV}$$



Downloaded from [1]

Magnetization in the quantum limit

Energy levels in B-field for different types of electrons [1]:

$$\epsilon_{n,k} = \begin{cases} \frac{\hbar e B}{m} (n + \gamma) + \frac{\hbar^2 k_z^2}{2m} & \text{Trivial } (\gamma = \frac{1}{2}) \\ \hbar v_F \sqrt{2B(n + \gamma) + k_z^2} & \text{Weyl } (\gamma = 0) \\ \hbar v_F \sqrt{2B(n + \gamma + C^2 \sin^2 \theta) + k_z^2} & \text{Dirac } (\gamma = 0) \end{cases}$$

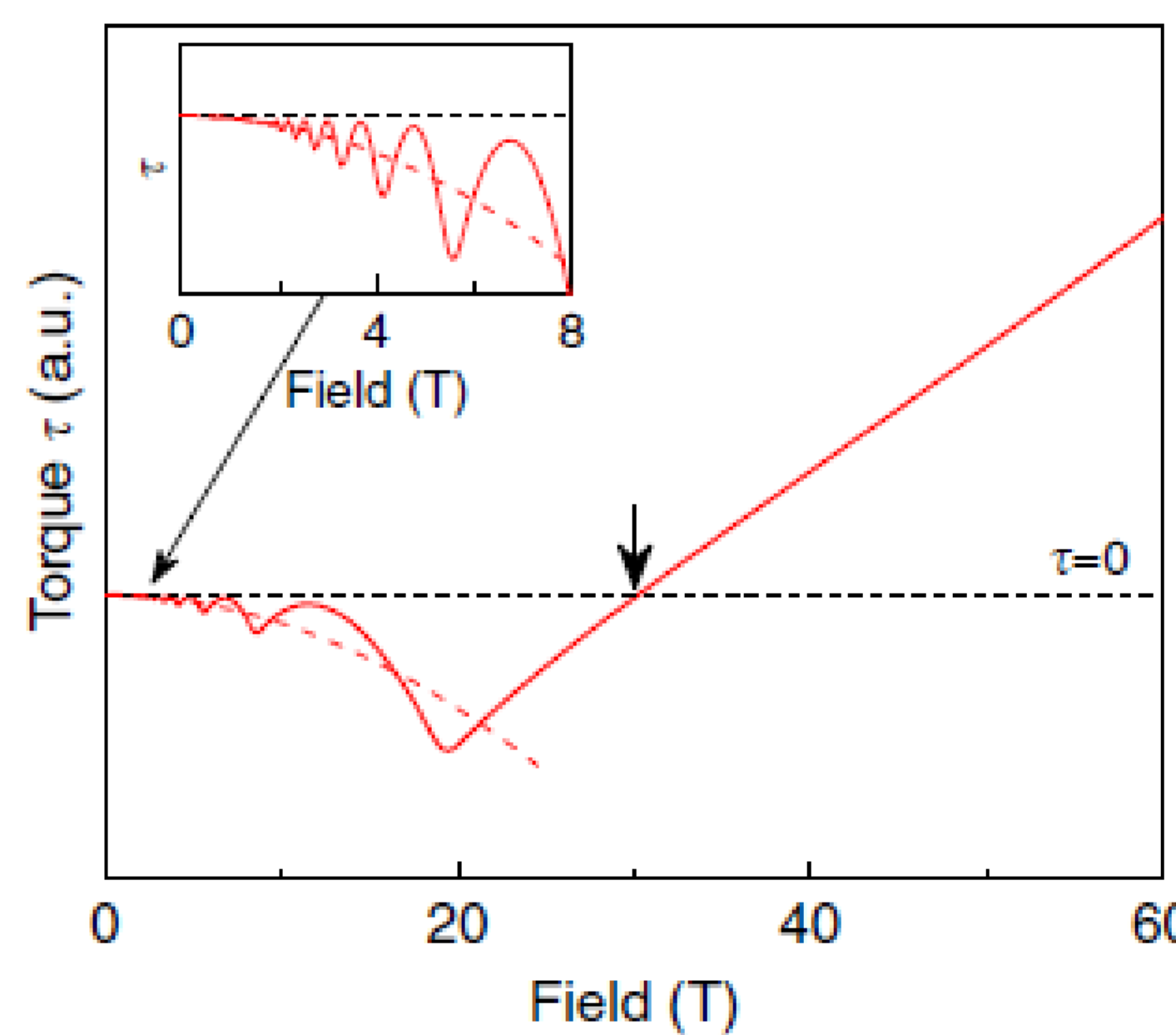
θ – angle between B and line connecting two Dirac points
C – material dependent parameter

Magnetization in quantum limit
 $M_{n=0} = -\frac{\partial \epsilon_{0,k}}{\partial B}$

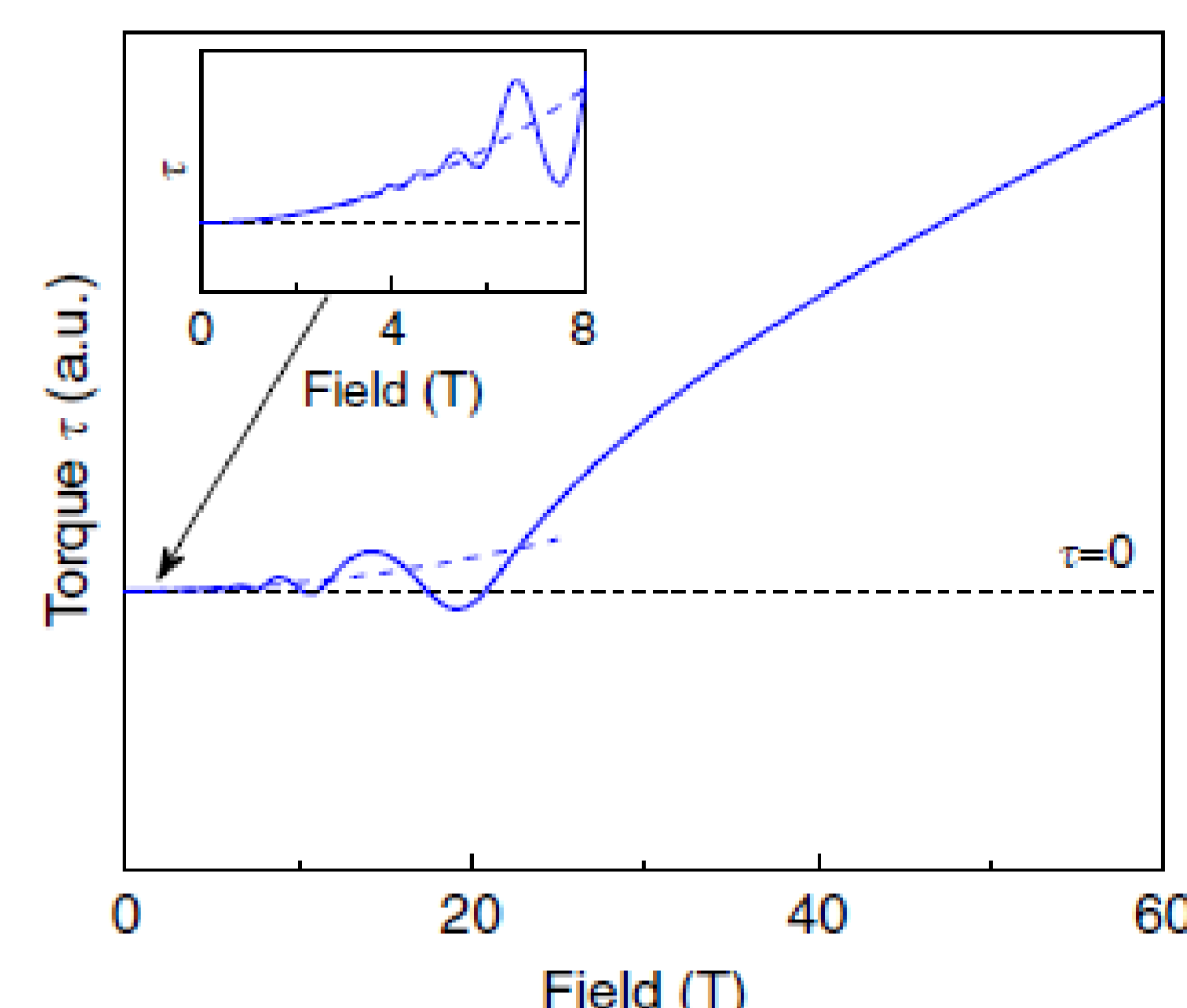
Dirac – B-field perpendicular to the symmetry axis \rightarrow massive fermion
– B-field in the direction of the symmetry axis \rightarrow weyl

Simulated behaviour of torque near the quantum limit [1]:

Downloaded from [1]



Weyl – a change from diamagnetic to paramagnetic response leading to the torque reversal.

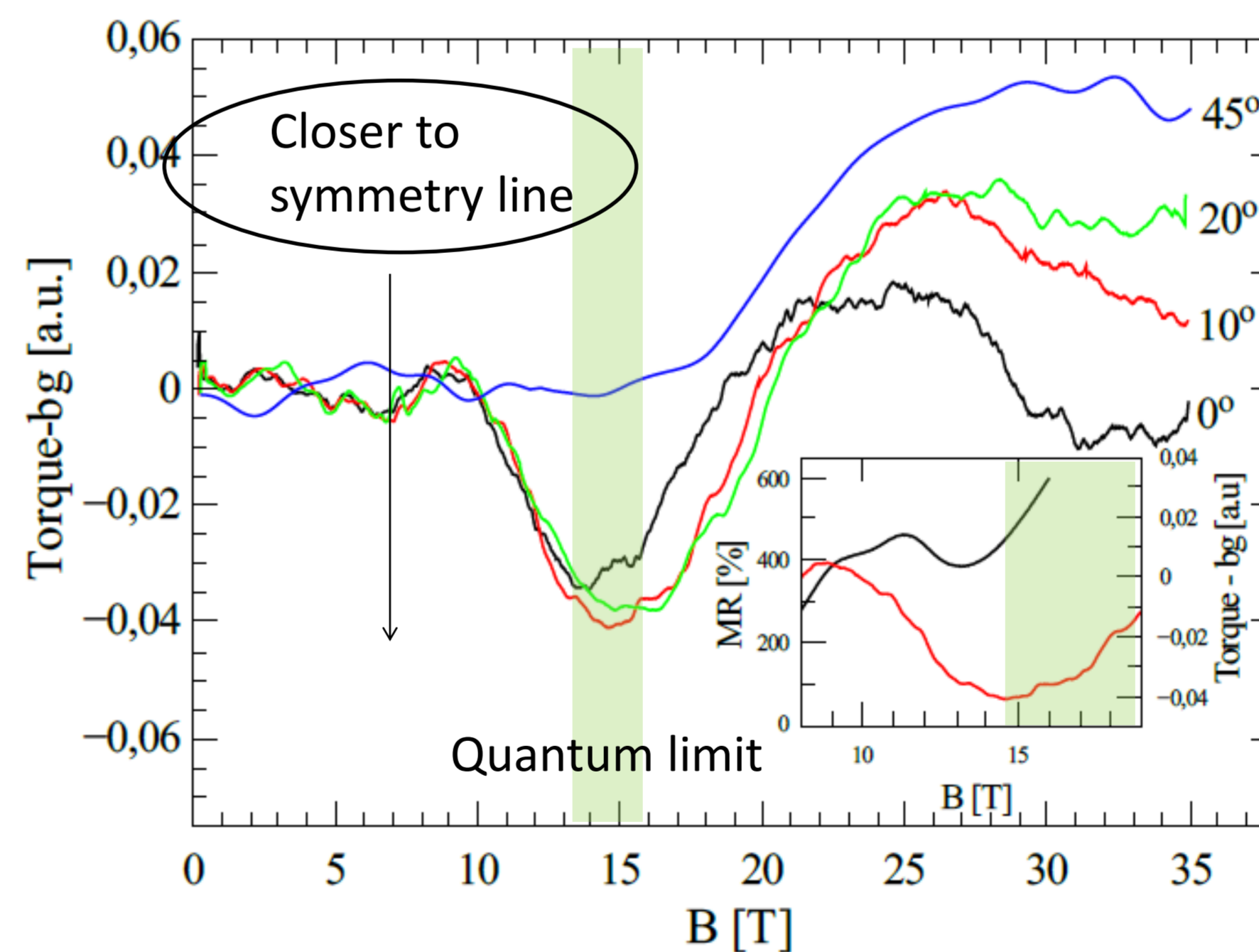
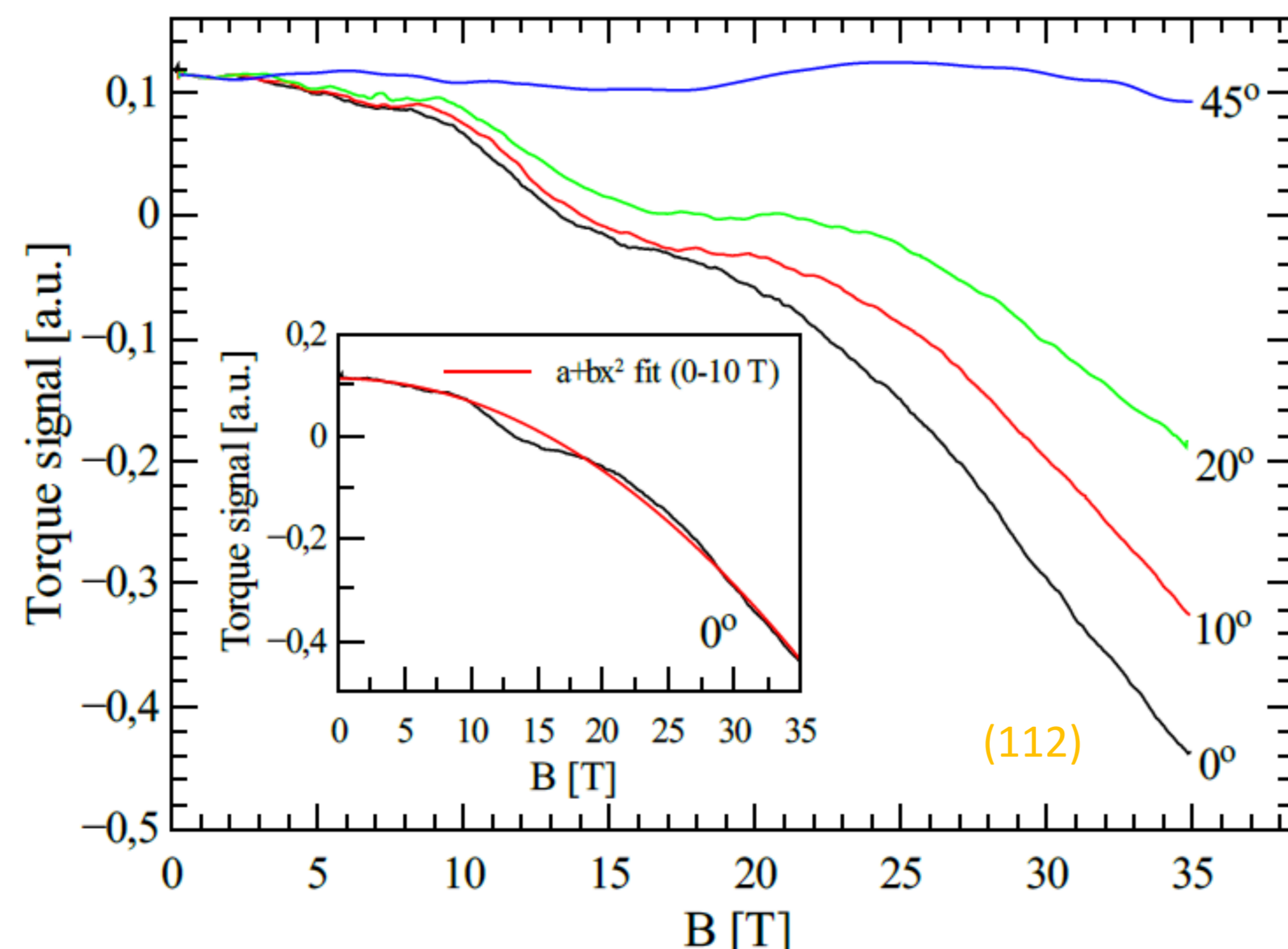
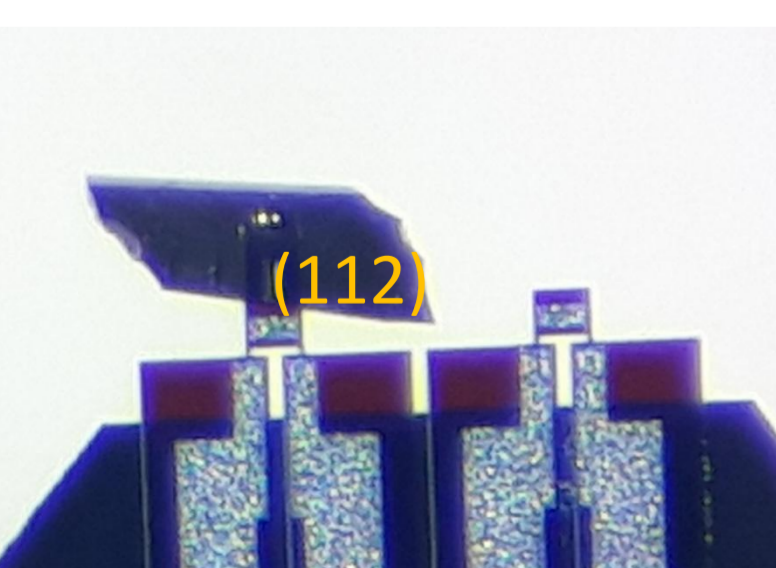
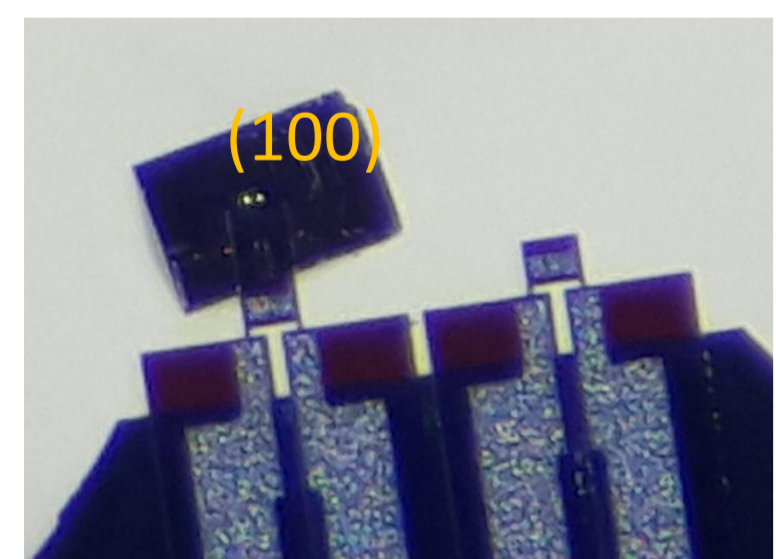


Massive fermion – no anomaly in magnetic response near quantum limit and no sign reversal in the torque.

Magnetic torque measurements

- By piezoresistive cantilever technique.
- Two different samples glued on different planes.
- At 4.2 K up to 35 T.

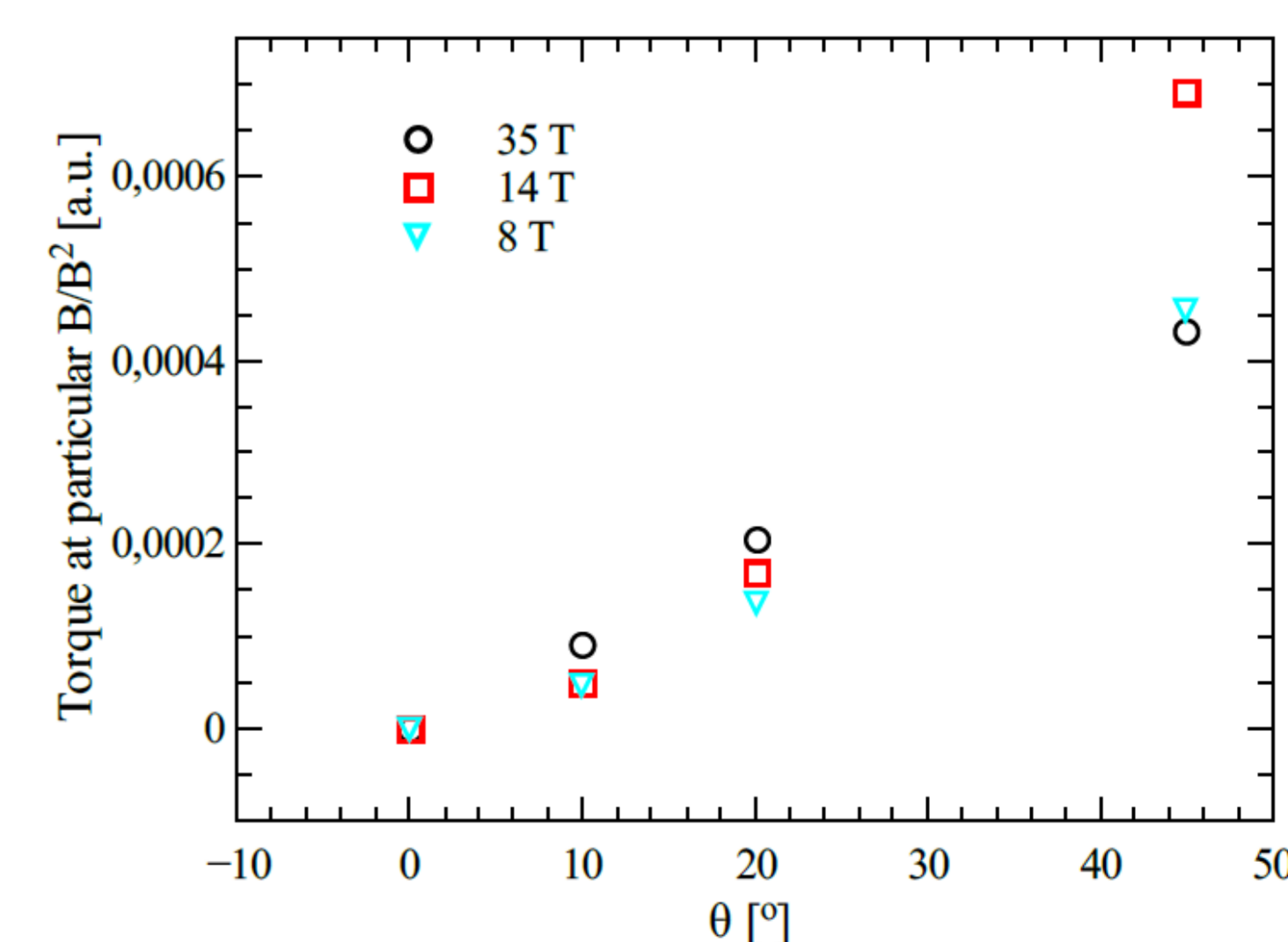
Angle is between field direction and normal to the sticking plane.



By entering the quantum limit the torque anomaly appears. Anomaly and change in torque slope is highly suppressed at 45°.

Quantum limit appears at the same field for different field directions indicating the sphericity of the Fermi surface (confirmed in SdH oscillations).

The same angular dependence below and above quantum limit at 45°.



Conclusion

- Successfully synthesized low charge concentration Cd_3As_2 samples with lower quantum limit (Cd_3As_2 usually has F around 45-50 T).
- From SdH oscillations Fermi surface is found to be spherical leading to very small torque signal. In samples with higher charge concentration the Fermi surface is ellipsoidal [2].
- Anomalous and angle dependent behavior in the torque near the quantum limit has been found.

References:

- [1] Moll, P. J. *et al.* (2016). Magnetic torque anomaly in the quantum limit of Weyl semimetals. *Nature Communications*, 7.
- [2] Borisenko, S. *et al.* (2014). Experimental realization of a three-dimensional Dirac semimetal. *Physical Review Letters*, 113(2).



Acknowledgments

This work has been fully supported by Croatian Science Foundation under the project No. 6216.